

How optics resolved the Einstein-Podolsky-Rosen paradox

F. Stef Roux

CSIR National Laser Centre

Presented at the

Department of Electrical, Electronic and Computer Engineering

University of Pretoria

22 October 2009



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Einstein-Bohr debate

Quantum mechanics:

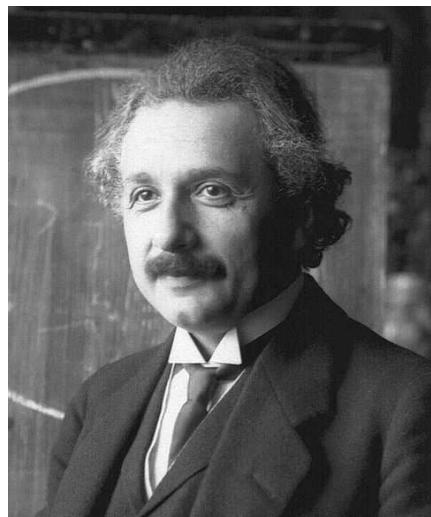
- ▷ Wave-particle duality
- ▷ Hilbert space
- ▷ Quantum superposition
- ▷ Heisenberg uncertainty

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$

Albert Einstein
and Neils Bohr
debating quantum
mechanics



Einstein-Podolsky-Rosen



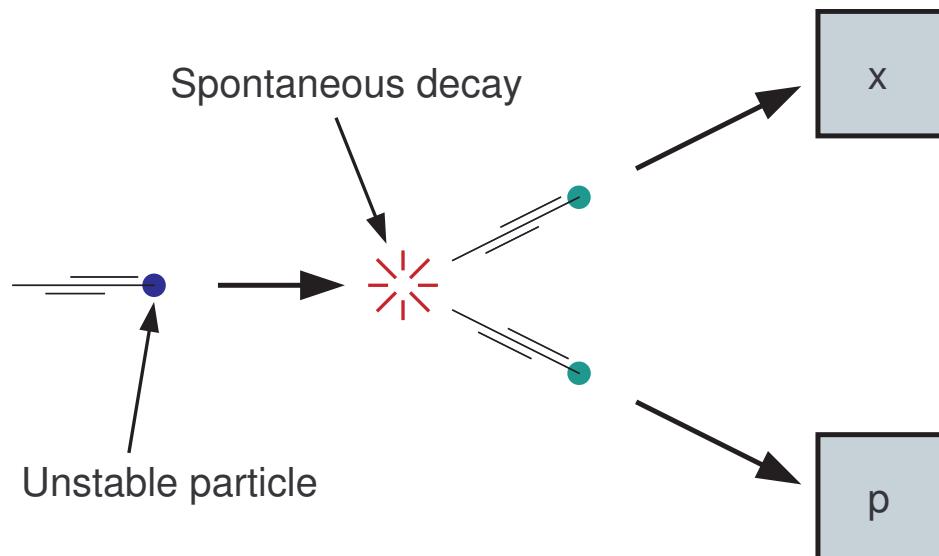
Albert Einstein



Boris Podolsky



Nathan Rosen



Quantum mechanics:
measurements on one
particle dictate the
state of the other particle.

Bell's inequality

Assumptions:

- ▷ Only local interactions
(no "spooky action at a distance")
- ▷ Unique reality
(not multiple realities)



Then:

John S. Bell

$$B(a, b, a', b') = |C(a, b) - C(a, b')| + |C(a', b) + C(a', b')| \leq 2$$

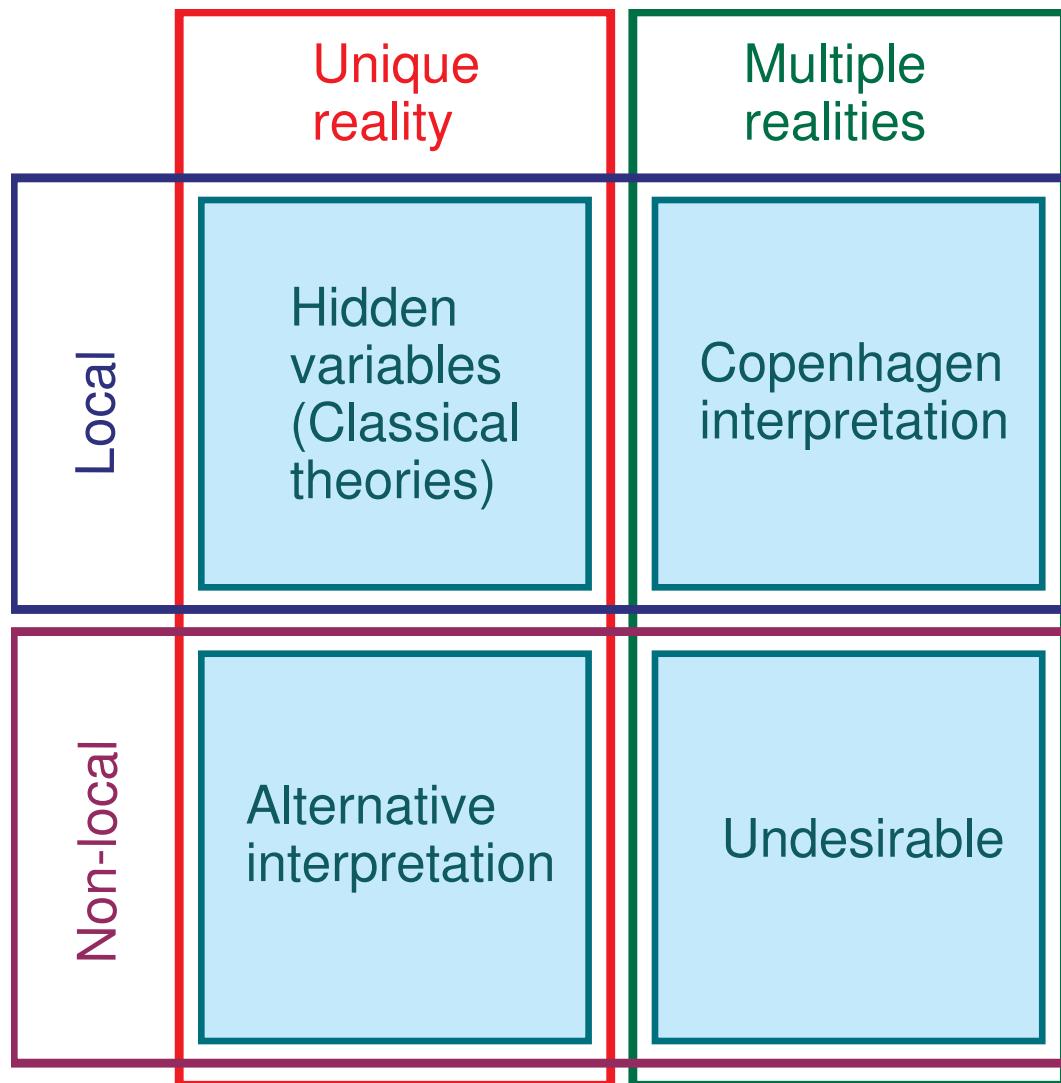
where C is the correlation between two measurements:

$$C(a, b) = E \{ A(a)B(b) \}$$

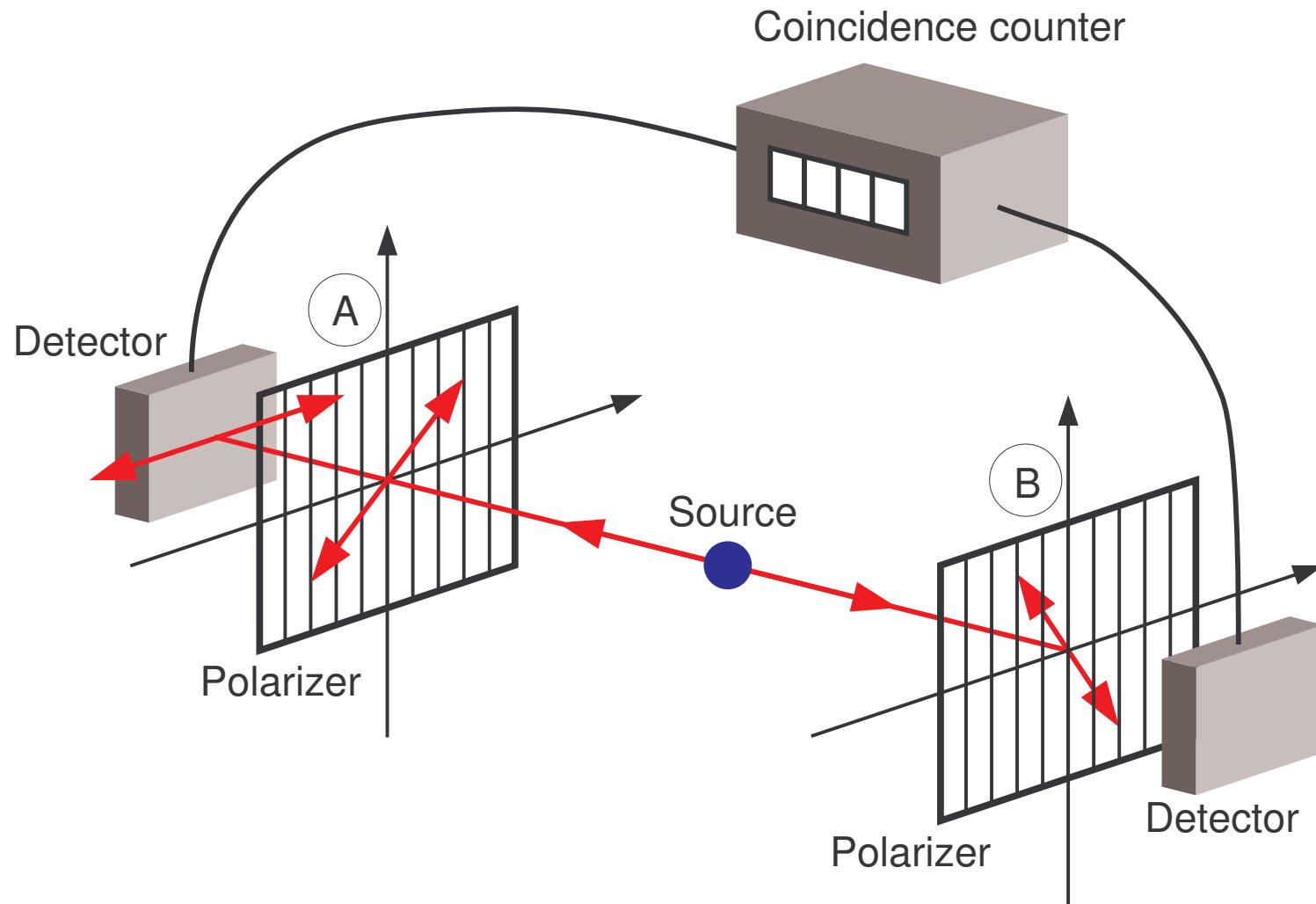


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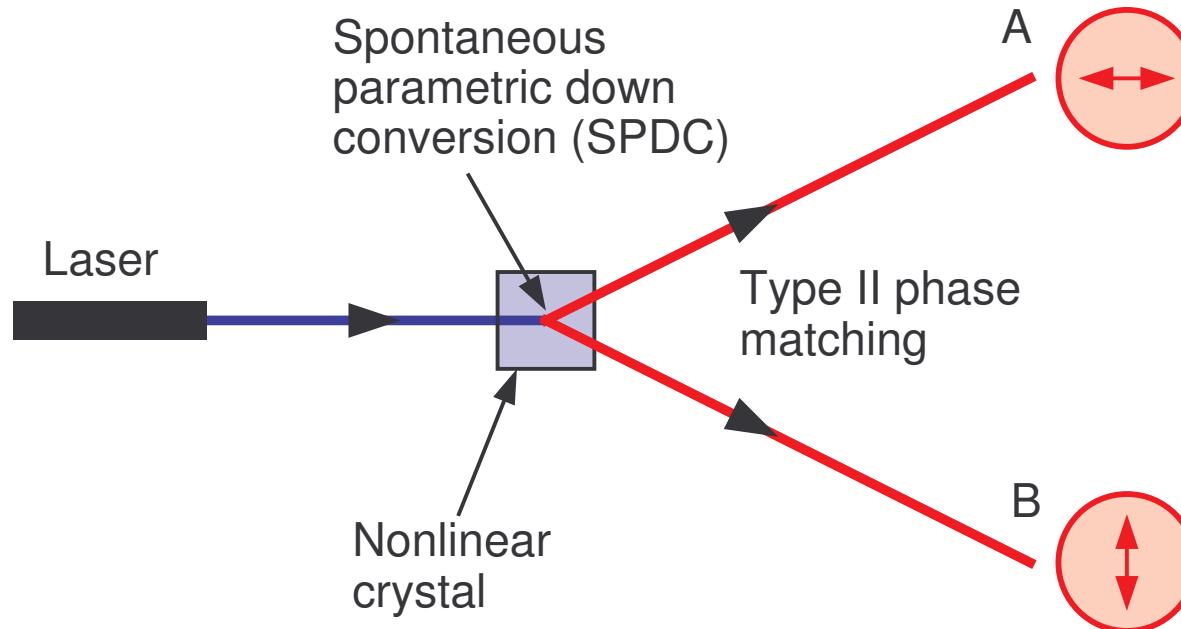
4 possibilities



Aspect experiment



Parametric down conversion



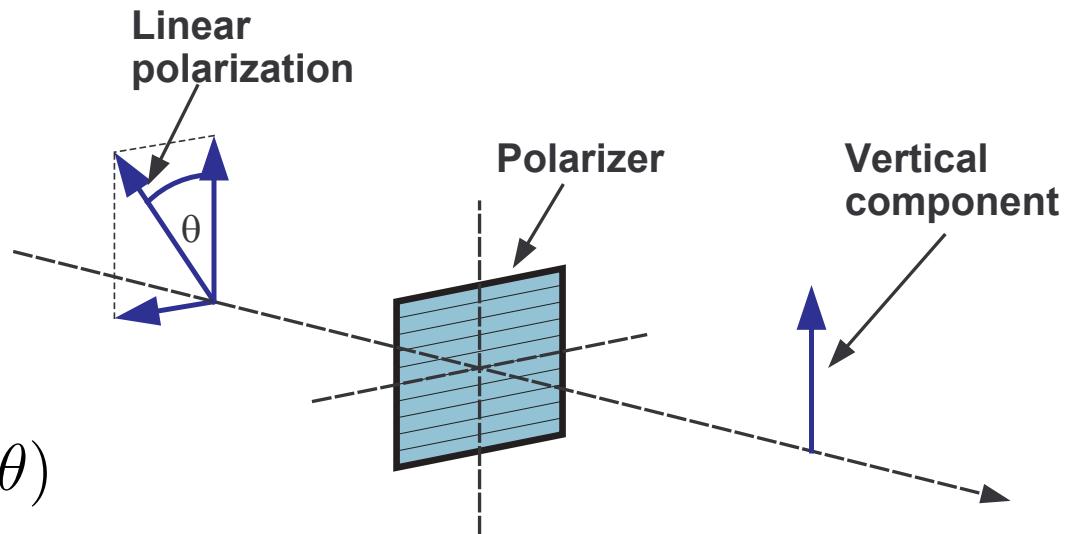
Type II phase matching \Rightarrow photons have perpendicular polarization: $\theta_B = \theta_A - \pi/2$

However, each beam is unpolarized
— contains all states of polarization.

Dichroic polarizers

Before considering the calculations, we first need to review dichroic polarizers.

$$\text{Intensity: } I_{ver} = I_{in} \cos^2(\theta)$$



where θ is the angle between vertical polarization and the input polarization.

$$\text{Probability to detect photon: } P(\theta) = \cos^2(\theta)$$

Classical calculations

Orientation angles of the two polarizers: α_A and α_B

Probability to detect photons at respective detectors:

$$P(A|\theta_A, \alpha_A) = \cos^2(\theta_A - \alpha_A)$$

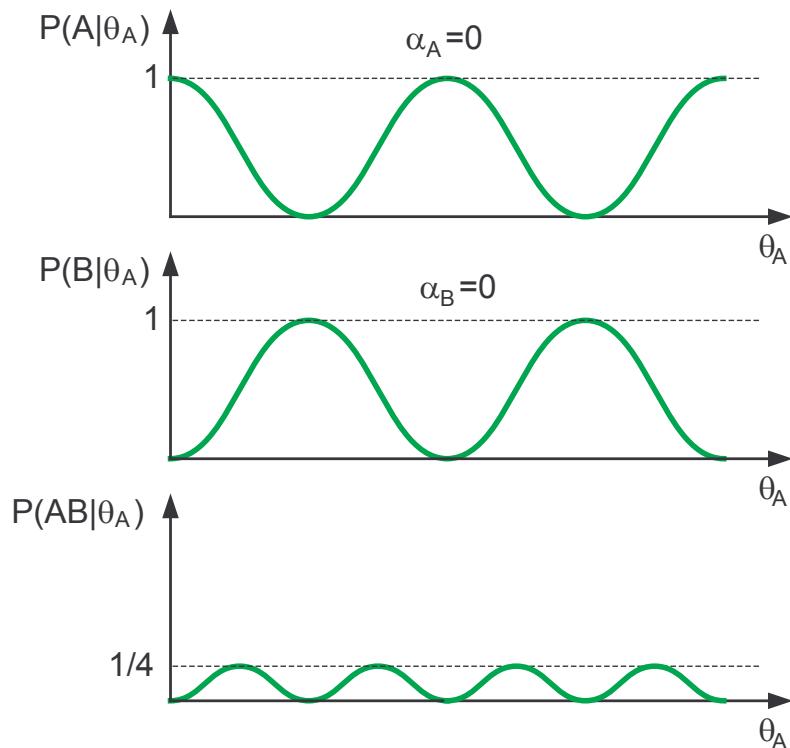
$$P(B|\theta_B, \alpha_B) = \cos^2(\theta_B - \alpha_B) = \sin^2(\theta_A - \alpha_B)$$

Detectors are far apart \rightarrow assume the probabilities of detecting photons are statistically independent

$$\Rightarrow P(AB|\theta_A, \alpha_A, \alpha_B) = P(A|\theta_A, \alpha_A)P(B|\theta_B, \alpha_B)$$

$$= \frac{1}{4} [\sin(\alpha_A - \alpha_B) + \sin(2\theta_A - \alpha_A - \alpha_B)]^2$$

Classical case



For the same orientation:
angle: ($\alpha_A = \alpha_B = \alpha$):

$$P(AB|\alpha, \alpha) = \frac{1}{8}$$

Integrate over all polarization states:

$$P(AB|\alpha_A, \alpha_B) = \frac{1}{2\pi} \int_0^{2\pi} P(AB|\theta_A, \alpha_A, \alpha_B) d\theta_A$$

$$= \frac{1}{8} + \frac{1}{4} \sin^2(\alpha_A - \alpha_B)$$

Quantum calculations

2-state maximally entangled bi-partite (2 photon) system:

$$|\psi\rangle = \frac{1}{\sqrt{2}} (|hor, ver\rangle - |ver, hor\rangle)$$

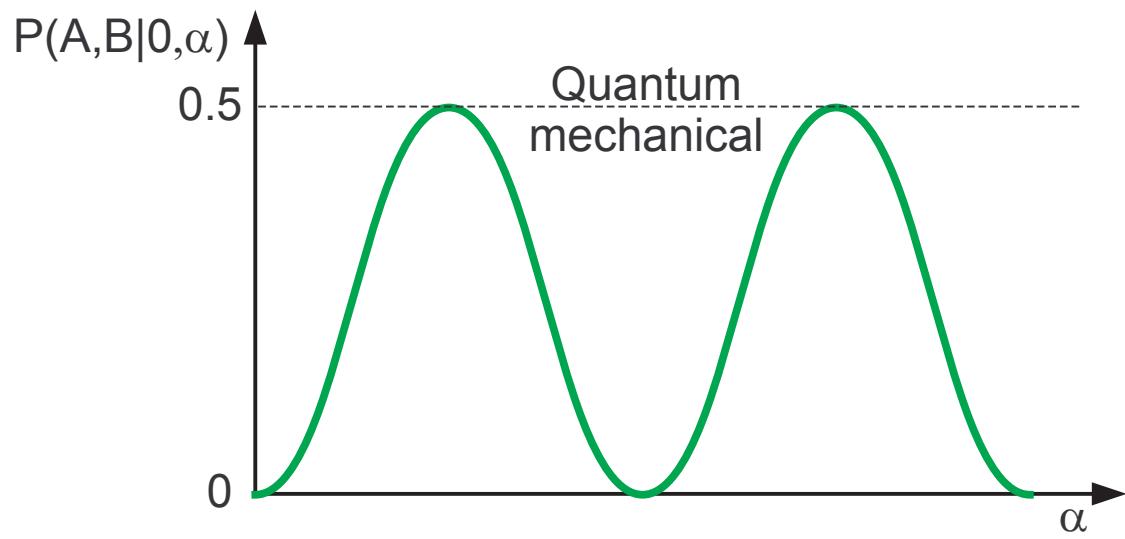
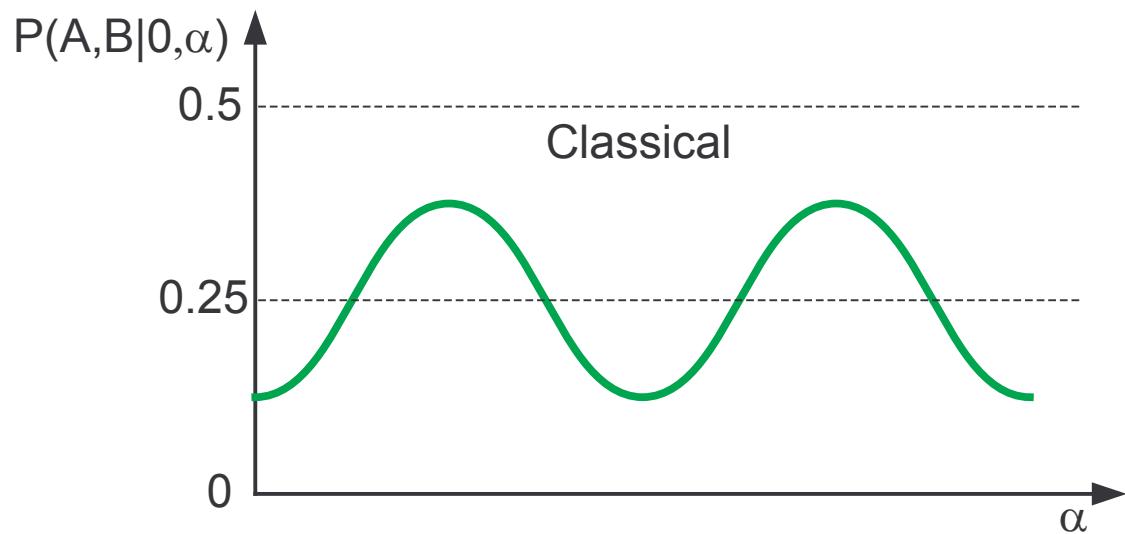
$$a_n = a_{n,hor} \cos \alpha_n + a_{n,ver} \sin \alpha_n \quad n = A, B$$

$$a_n^\dagger = {a_{n,hor}}^\dagger \cos \alpha_n + {a_{n,ver}}^\dagger \sin \alpha_n \quad n = A, B$$

$$P(A, B | \alpha_A, \alpha_B) = \langle \psi | a_A^\dagger a_B^\dagger a_A a_B | \psi \rangle = \frac{1}{2} \sin^2(\alpha_A - \alpha_B)$$

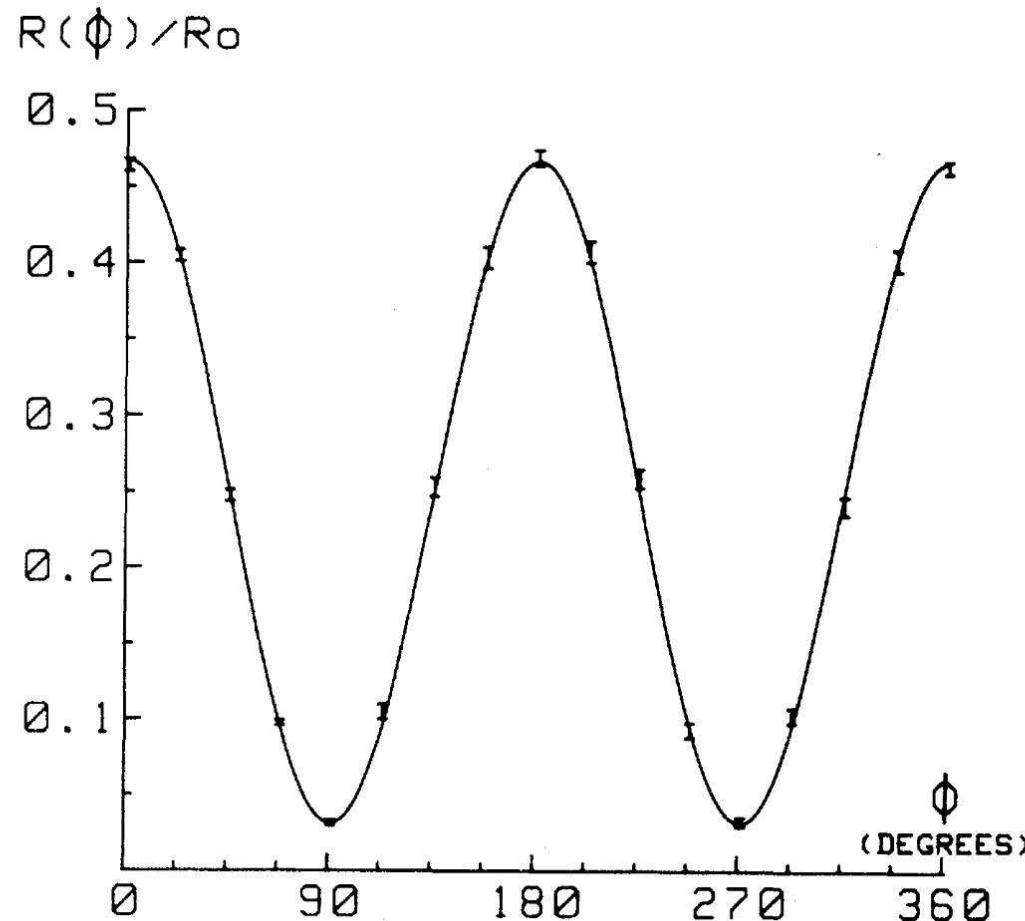
So for $\alpha_A - \alpha_B = 0$ (same orientation) $P(A, B | \alpha, \alpha) = 0$

Comparison



Experimental results

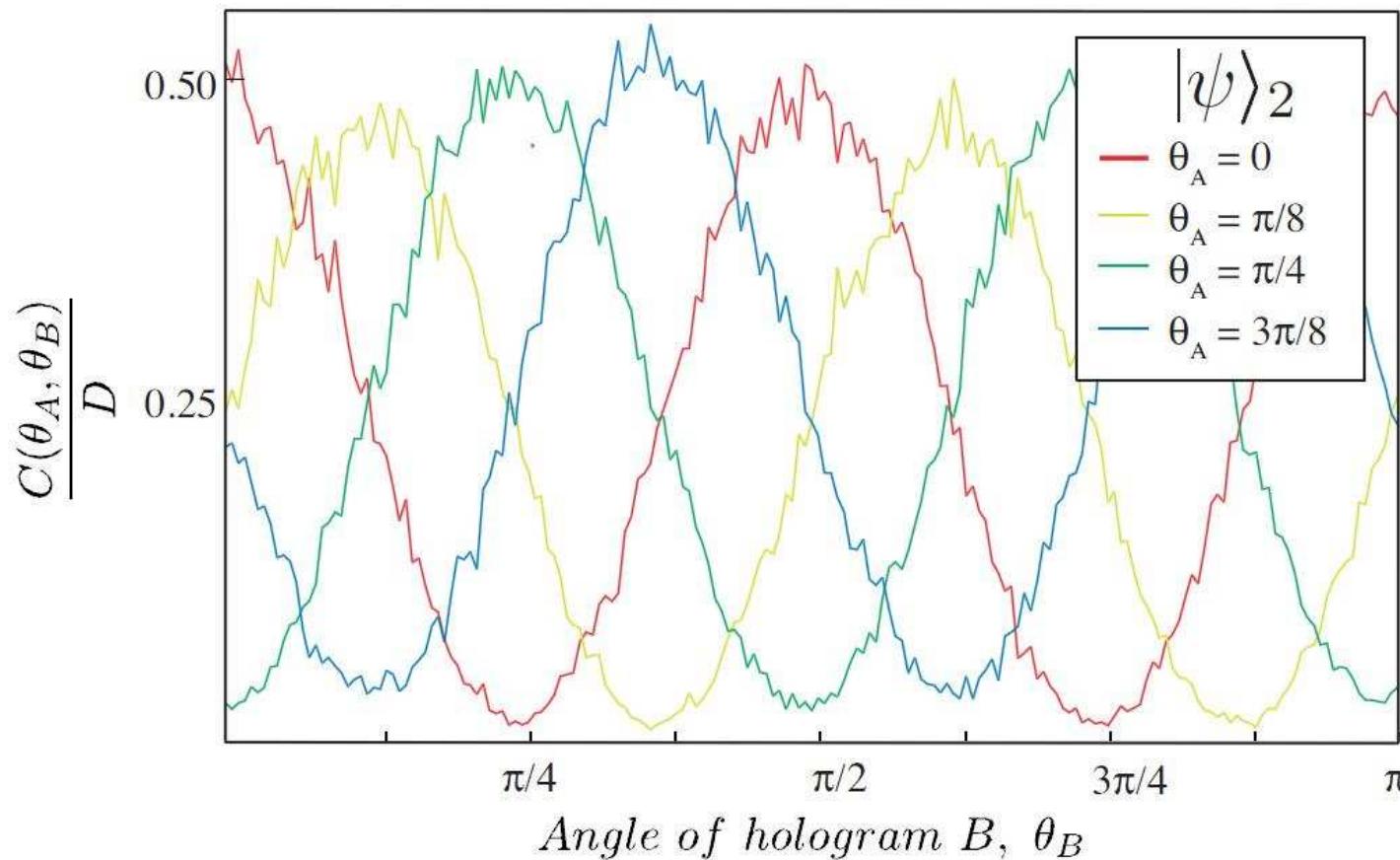
For polarization ... ^a



^a A. Aspect, et al., Phys. Rev. Lett., Vol 47, p. 460 (1981).

More experimental results

(... and for orbital angular momentum) ^a



^a J. Leach, et al., Optics Express, Vol 17, p. 8287 (2009).

Test of Bell's inequality

Theoretically:

Classical case: $C_C(\alpha_A, \alpha_B) = -\frac{1}{2} \cos(2\alpha_A - 2\alpha_B)$

Quantum case: $C_Q(\alpha_A, \alpha_B) = -\cos(2\alpha_A - 2\alpha_B)$

Maximum violation for:

$$\alpha_A = 0, \alpha_B = \frac{3\pi}{8}, \alpha_A' = -\frac{\pi}{4}, \alpha_B' = \frac{\pi}{8}$$

Gives

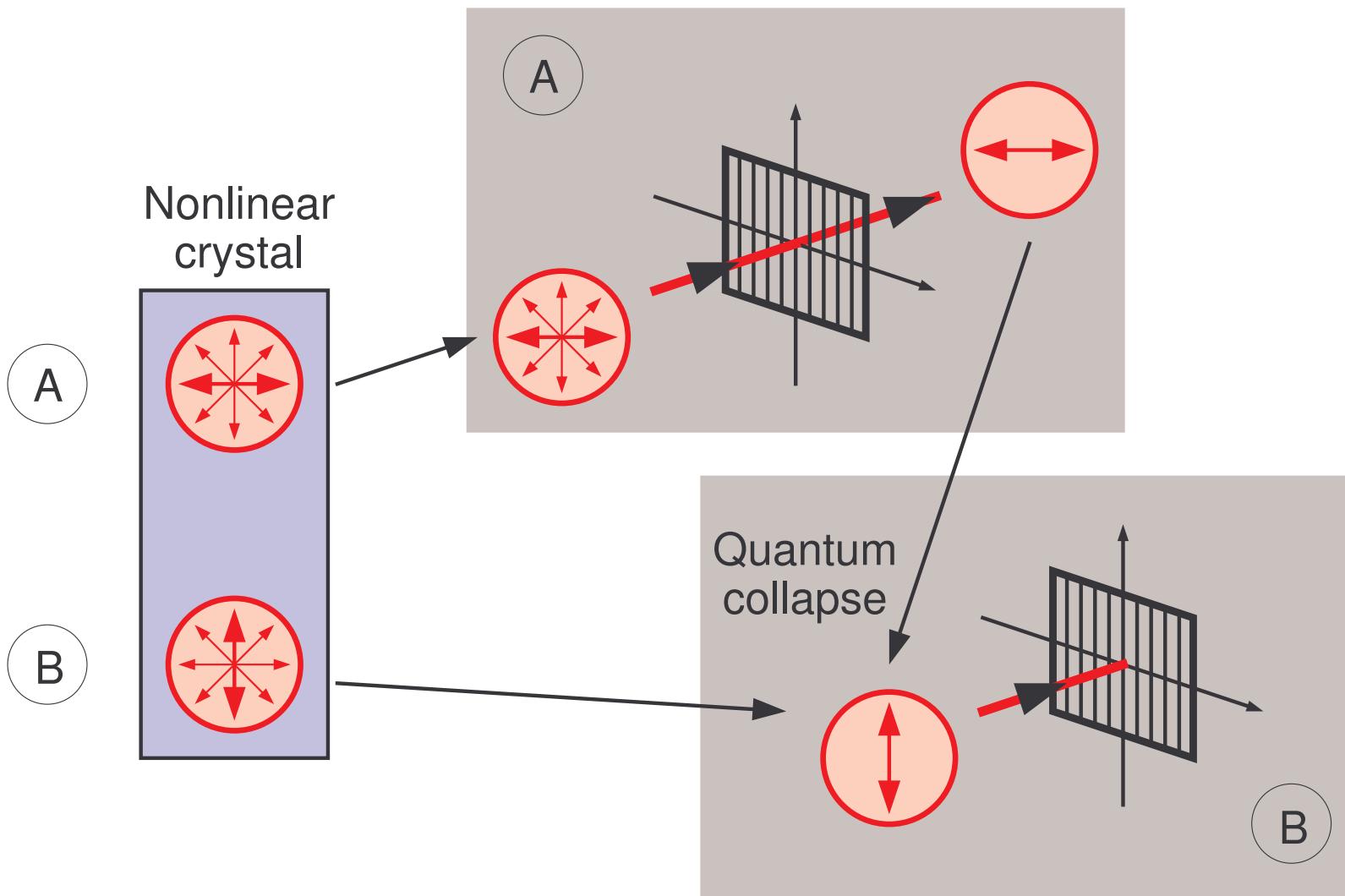
$$B_C(\alpha_A, \alpha_B, \alpha_A', \alpha_B') = 1.4 < 2$$

and

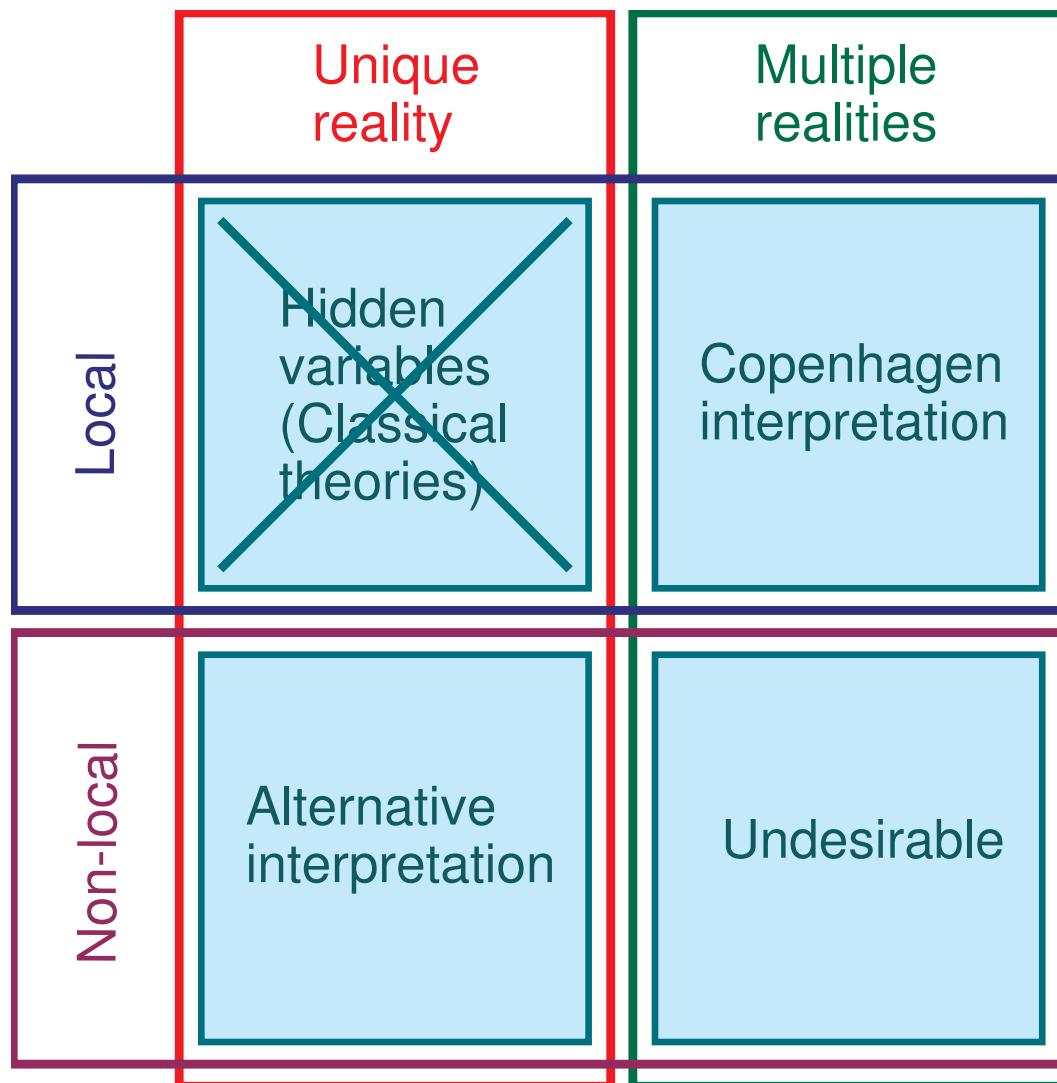
$$B_Q(\alpha_A, \alpha_B, \alpha_A', \alpha_B') = 2.8 > 2$$

Experimentally Bell's (or related) inequality was violated by several standard deviations.

How does it work?



What does it mean?



Where can we use it?

- ▷ Quantum communication
 - quantum cryptography
- ▷ Quantum computing
 - efficient factorization
- ▷ Ghost imaging
- ▷ Exotic photon sources

